PART V

Stellar Feedback

1. What Escapes From Stars

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Energy: E_{\text{tot}} = E_{\text{SNII}} + E_{\text{SNIa}} + E_{\text{wind}}

Momentum: p_{\text{tot}} = p_{\text{SNII}} + p_{\text{wind}} + p_{\text{rad}}

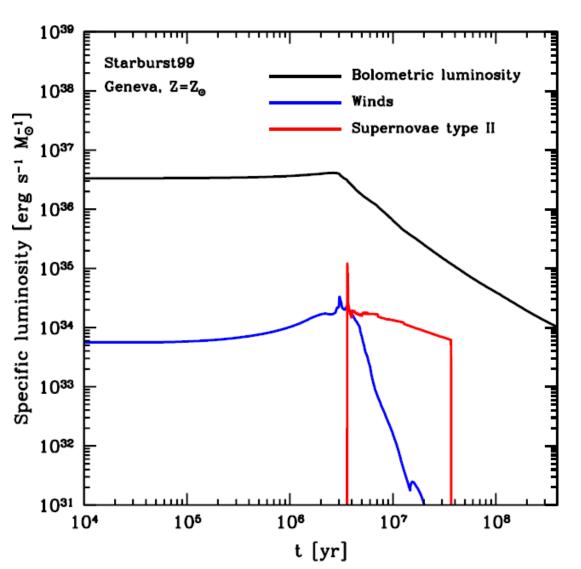
Mass loss: m_{\text{tot}} = m_{\text{SNII}} + m_{\text{SNIa}} + m_{\text{wind}} + m_{\text{loss}}

Metals: m_{Z,\text{tot}} = m_{Z,\text{SNII}} + m_{Z,\text{SNIa}} + m_{Z,\text{wind}} + m_{Z,\text{loss}}
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(+ dust)

(+ cosmic rays)

Two Things That Matter: Energetics And Timing

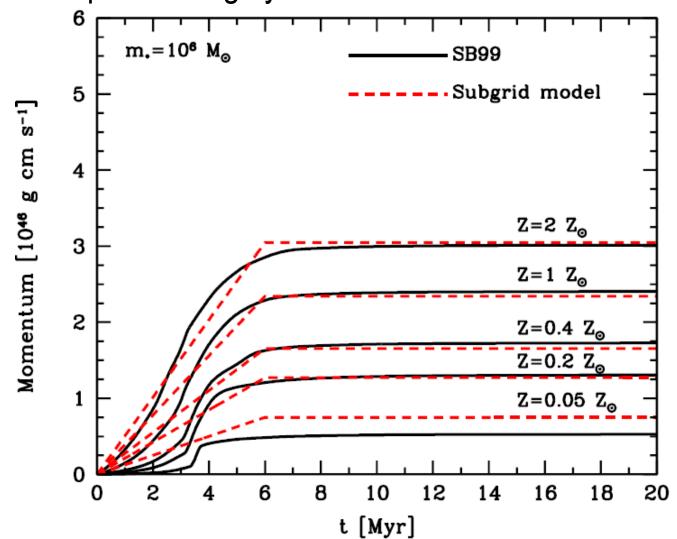


Important fact:
 the largest (by far)
 energy output of
 stars is in light!

 Hands-on exercise: add other important lines on this plot.

Stellar Winds

• Wind power roughly scales as $Z^{0.4}$.



Radiation Pressure

 For heavily obscured stars most of their light is absorbed; since photons have momentum, absorbing all light from a star/star cluster of luminosity L injects momentum into the surrounding gas:

$$\dot{p}_1 = \frac{L}{c}$$

 However, energy is conserved – the absorbed energy is irradiated by dust in the IR. If there is enough dust, it will be optically thick to its own IR radiation. That IR radiation does work on (=injects extra momentum into) the gas.

$$\dot{p}_{\text{tot}} = \left(1 - e^{-\tau_{\text{UV}}} + \tau_{\text{IR}}\right) \frac{L}{c}$$

The τ_{IR} Term: Calculus

• The extra $au_{
m IR}$ term can be derived easily: since energy is conserved, the radiation flux at each radius R is still

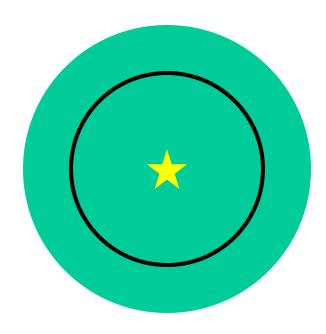
$$F_R = \frac{L}{4\pi R^2}$$

• Hence, the momentum imparted between ${\cal R}$ and

$$R + dR$$
 is

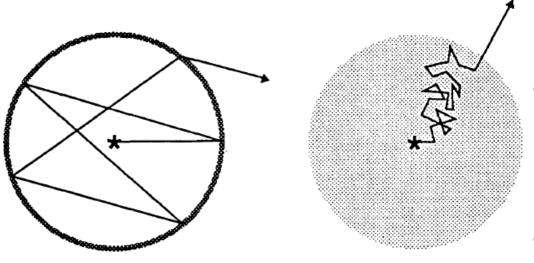
$$d\dot{p}_{\rm IR} = 4\pi R^2 \frac{F_R}{c} \kappa dR = \frac{L}{c} d\tau$$

$$\dot{p}_{\rm IR} = \tau_{\rm IR} \frac{L}{c}$$

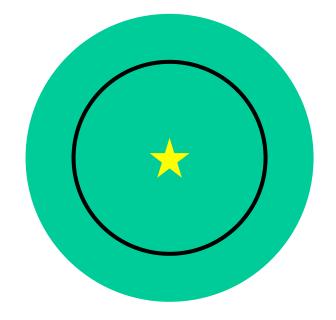




The τ_{IR} Term: Geometry



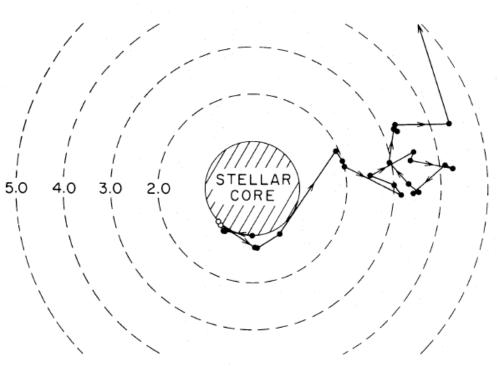
- For a thin shell of optical depth \(\tau \) there is a clear geometric explanation.
- What about a real case (continuous medium)?



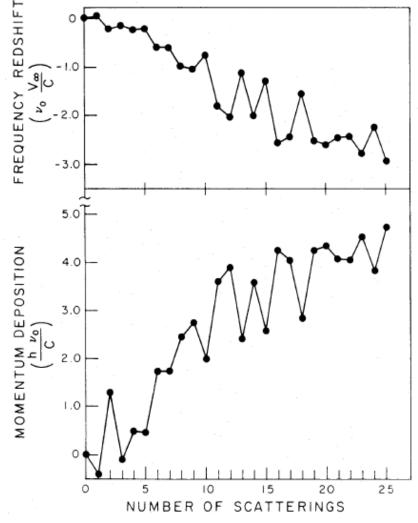
 Recall: "ogres are like onions". That applies to spherical clouds too.

The τ_{IR} Term In Action

An example of photon scattering:



• Note: number of scatterings is τ^2 , not τ .



Radiation Pressure

• What about $\tau_{\rm IR}$?

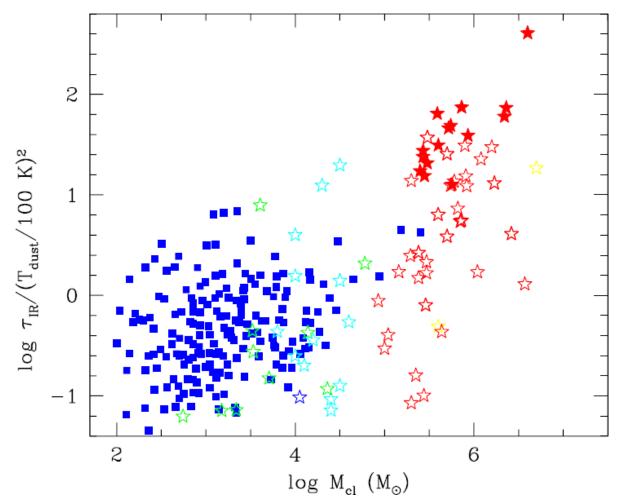
$$\tau_{\rm IR} = \varkappa_{\rm IR} \Sigma_g$$

$$\varkappa_{\rm IR} \approx 3 \frac{\rm cm^2}{\rm g} \left(\frac{T_d}{100 \rm K} \right)^2$$

(Hot dust absorbs way more).

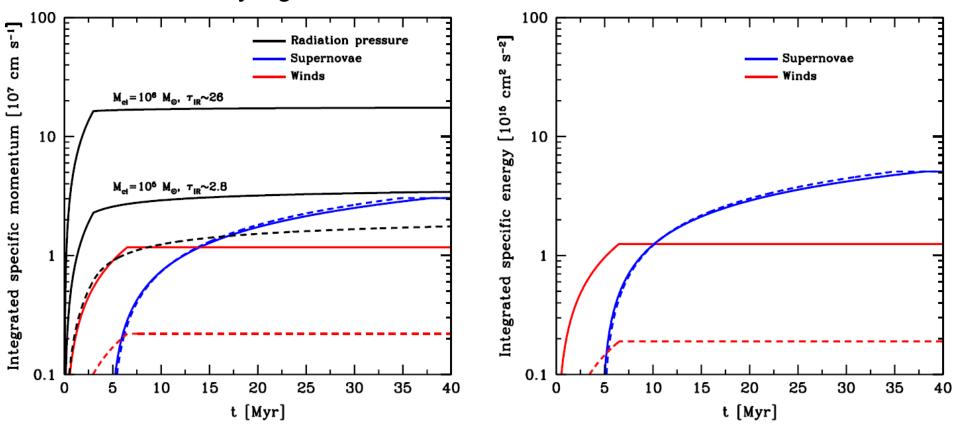
Radiation Pressure

 Radiation pressure is particularly important for large stellar clusters.

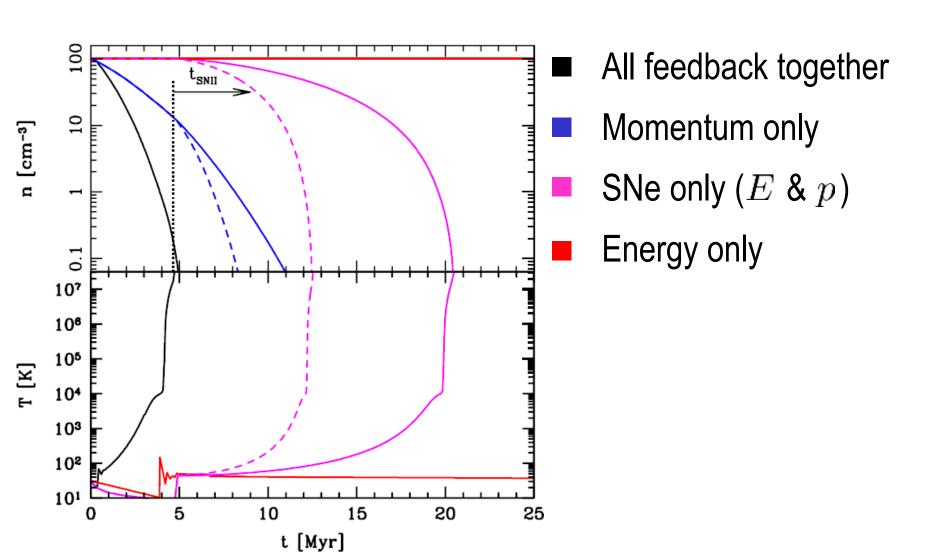


All Feedback Together

- Important lesson: supernovae go off after most of other feedback already did its share.
- For destroying molecular clouds SNe are irrelevant.



Now, Let's Simulate!

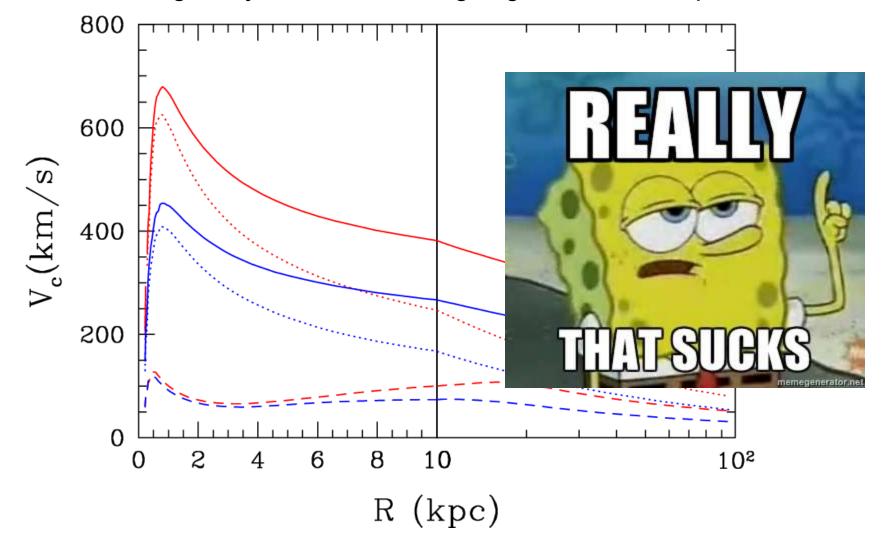


2. Unconventional Marriage: Feedback & SF

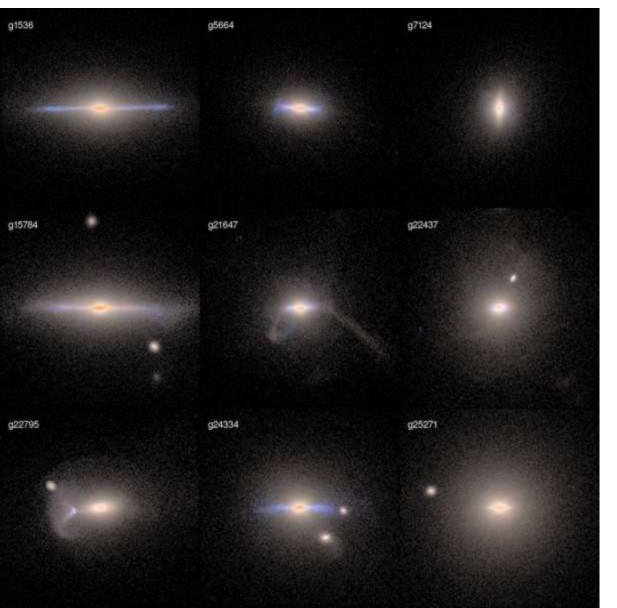
- Recall: $\tau_{\rm SF}$ is long (~ 1 Gyr), while molecular clouds are short lived (10-20 Myr).
- Hence, star formation is inefficient only a small fraction of all molecular gas goes into stars in a molecular cloud during its lifetime.
- A natural conclusion: since feedback comes from stars, it should be inefficient too.

How About This?

The same galaxy that has the right gaseous halo profile...



Simulated Galaxies...

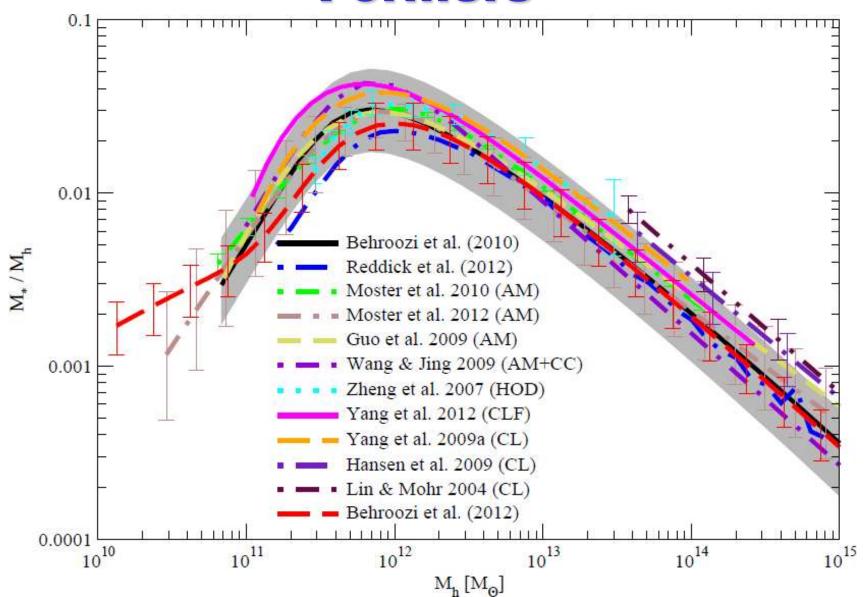


- MUGS galaxies (Stinson et al 2010)
- All simulated galaxies have large bulges.

...And Real Ones



Galaxies Are Inefficient Star-Formers



SF-Feedback Dilemma

 The challenge in understanding star formation and feedback is that star formation is inefficient, but feedback must be efficient in expelling a significant fraction of baryons from MW-like galaxies and even larger fractions for low and higher mass galaxies.

3. Miracles Do Happen

 Historically, there were many attempts to model feedback in cosmological and galactic simulations with sub-grid models.

None of them was particular successful...

Yepes et al 1997

$$\left(\frac{\mathrm{d}\rho_*}{\mathrm{d}t}\right)_{\mathrm{star\ formation}} = \frac{(1-\beta)\rho_{\mathrm{c}}}{t_*}.$$

$$\left(\frac{\mathrm{d}\rho_{\mathrm{h}}}{\mathrm{d}t}\right)_{\mathrm{star\ formation}} = \frac{\beta\rho_{\mathrm{c}}}{t_{*}}.$$

$$\left(\frac{\mathrm{d}\rho_*}{\mathrm{d}t}\right)_{\mathrm{star\ formation}} = \frac{(1-\beta)\rho_{\mathrm{c}}}{t_*}. \qquad \left(\frac{\mathrm{d}\rho_{\mathrm{h}}}{\mathrm{d}t}\right)_{\mathrm{evap}} = -\left(\frac{\mathrm{d}\rho_{\mathrm{c}}}{\mathrm{d}t}\right)_{\mathrm{evap}} = \frac{A\beta\rho_{\mathrm{c}}}{t_*}$$

Correspondingly, the net energ phase by supernovae is

$$\left(\frac{\mathrm{d}\rho_{\mathrm{h}}\varepsilon_{\mathrm{h}}}{\mathrm{d}t}\right)_{\mathrm{SN}} = \frac{\beta\rho_{\mathrm{c}}}{t_{*}} \left[\varepsilon_{\mathrm{SN}} + (A+1)\varepsilon_{\mathrm{c}}\right].$$

Gnedin 1998

$$\frac{\mathrm{d}\rho_*}{\mathrm{d}t} = \left(\frac{\mathrm{d}\rho_*}{\mathrm{d}t}\right)_0 (1 - f_{\text{hot}}),$$

$$v = q \frac{4\pi}{3} R_{\rm S}^3 t_{\rm S},$$

$$R_{\rm S} = 25.7 \text{ pc} \frac{E_{51}^{31/98}}{\xi^{5/98} n_0^{18/49} v_{100}^{3/7}},$$

$$t_{\rm S} = 7.57 \times 10^4 \text{ yr } \frac{E_{51}^{31/98}}{\xi^{5/98} n_0^{18/49} v_{100}^{10/7}}.$$

$$f_{\text{hot}} = 1 - \exp(-Q),$$

where (McKee 1990; Silk 1997)

$$Q = v \frac{1}{m_{\rm SN}} \frac{\mathrm{d}\rho_*}{\mathrm{d}t}.$$

Springel & Hernquist 2003

$$\frac{\mathrm{d}\rho_{\star}}{\mathrm{d}t} = \frac{\rho_{\mathrm{c}}}{t_{\star}} - \beta \frac{\rho_{\mathrm{c}}}{t_{\star}} = (1 - \beta) \frac{\rho_{\mathrm{c}}}{t_{\star}}.$$

$$\frac{\mathrm{d}\rho_{\mathrm{c}}}{\mathrm{d}t} = -\frac{\rho_{\mathrm{c}}}{t_{\star}} - A\beta \frac{\rho_{\mathrm{c}}}{t_{\star}} + \frac{1 - f}{u_{\mathrm{h}} - u_{\mathrm{c}}} \Lambda_{\mathrm{net}}(\rho_{\mathrm{h}}, u_{\mathrm{h}}),$$

$$\frac{\mathrm{d}\rho_{\mathrm{h}}}{\mathrm{d}t} = \beta \frac{\rho_{\mathrm{c}}}{t_{\star}} + A\beta \frac{\rho_{\mathrm{c}}}{t_{\star}} - \frac{1 - f}{u_{\mathrm{h}} - u_{\mathrm{c}}} \Lambda_{\mathrm{net}}(\rho_{\mathrm{h}}, u_{\mathrm{h}}).$$

Teyssier et al 2013

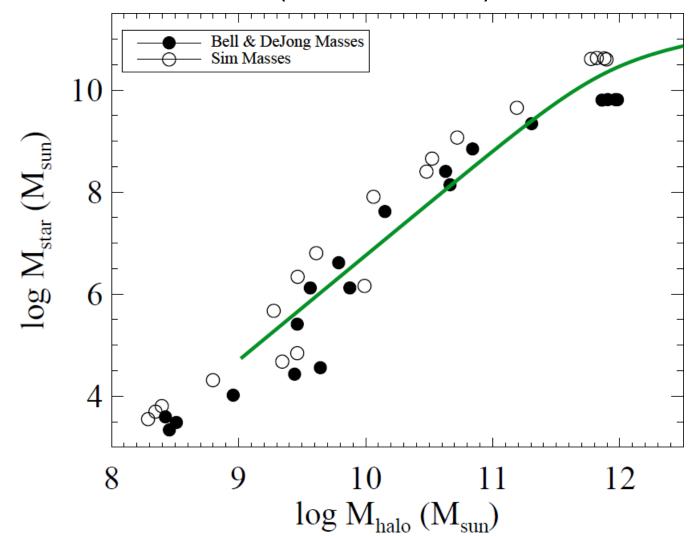
$$\dot{\rho}_* = \epsilon_* \frac{\rho_{\rm gas}}{t_{\rm ff}} \quad {\rm if} \quad \rho_{\rm gas} > \rho_*.$$

$$\rho \frac{D\epsilon_{turb}}{Dt} = \dot{E}_{inj} - \frac{\rho\epsilon_{turb}}{t_{diss}}$$

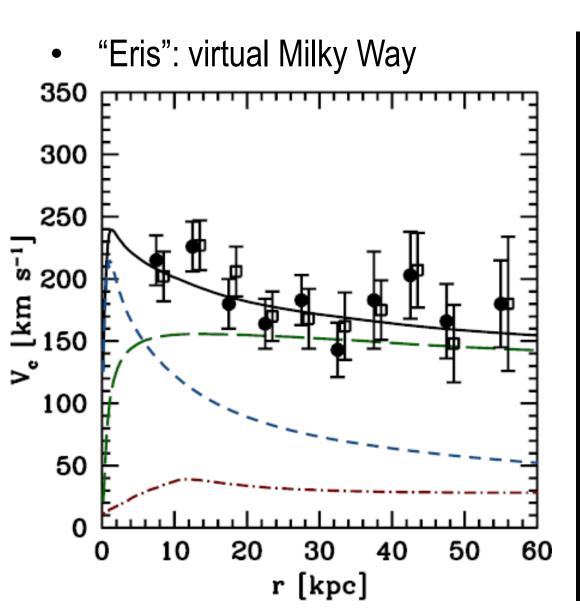
$$\rho \frac{D\epsilon_{\text{thermal}}}{Dt} = \dot{E}_{\text{inj}} - P_{\text{thermal}} \nabla \cdot \mathbf{v} - n_H^2 \Lambda$$

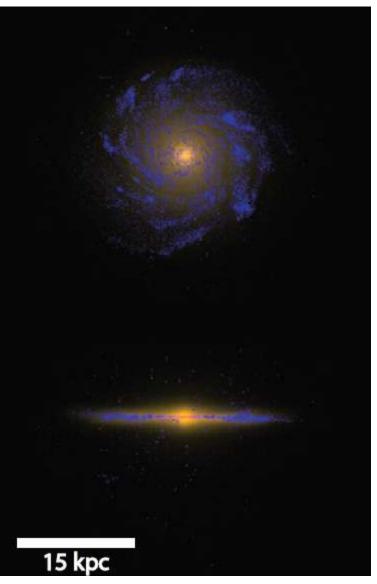
The World Is Full Of Magic

Munshi et al 2012 (Gasoline team):



The World Is Full Of Magic





And The Magic Is...

- All one needs to do to make a realistic galaxy in a cosmological simulation is to switch off cooling in a star forming region for ~10 Myr.
- Things to note:
 - As such, it is a trick, it has no physical justification.
 - The fact that it works may be a pure coincidence.
 - Alternatively, it can be a manifestation of a real physical process that operates on <1pc scales, but its consequences on ~100pc appear as if cooling was switched off.

Nature Of Delayed Cooling

Quiz: Which small-scale physical process behaves as delayed cooling on large scales:

- A. Radiation pressure massive (=bright) stars live for ~10 Myr.
- B. Thermal pressure of coronal gas cooling time of coronal gas ~10 Myr.
- C. Kinetic energy of internal motions inside molecular clouds molecular cloud lifetimes are ~10 Myr.
- D. Cosmic Ray pressure cosmic ray energy loss time in molecular clouds ~10 Myr.

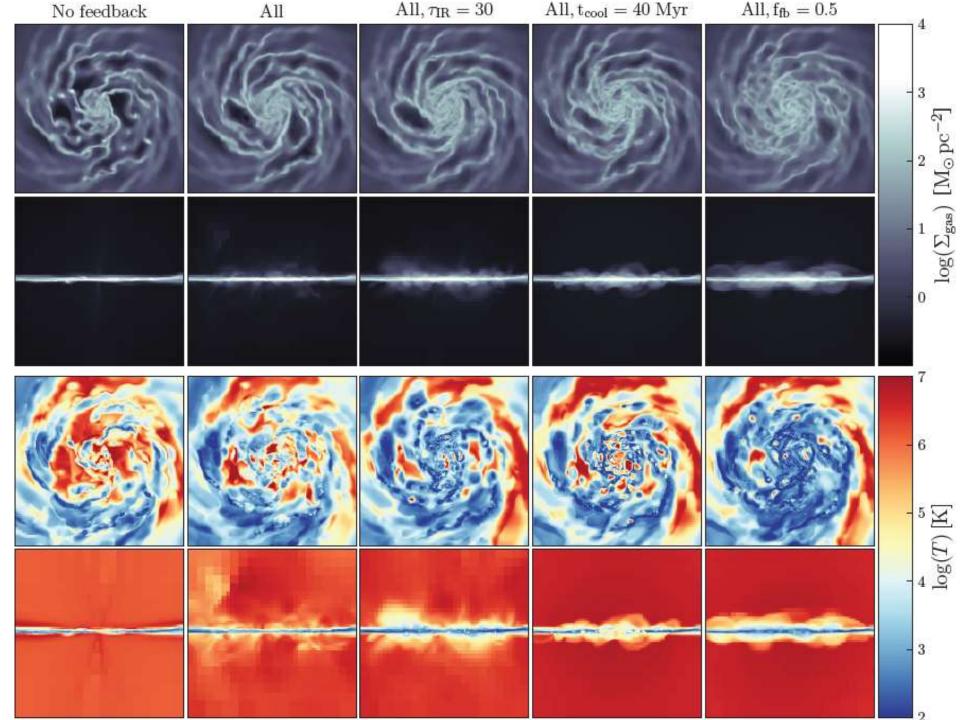
4. Star Formation Model And Feedback

- The efficiency of the feedback does depend on your adopted star formation model:
- The canonical $au_{\rm SF} = au_{\rm ff}/\epsilon_{\rm ff}$ model $(\dot{
 ho}_* \propto
 ho_{\rm mg}^{3/2})$ converts gas into stars too fast at early times, and requires extra "early" feedback on top of delayed cooling.
- The modern $au_{\rm SF}={
 m const}$ model $(\dot{
 ho}_*\propto
 ho_{
 m mg})$ avoids the need for early feedback, but does not concentrate feedback enough.
- Looks like existing SF models actually ran out of steam...

5. Towards The Future

 We need SF folks to tell us how/which/when molecular gas makes stars – looks like not all gas makes stars equally well.

- Different feedback modes do produce significant differences in the details of modeled galaxies.
- The challenge now is to figure out which one(s) of several possible physical mechanisms is/are the primary stellar feedback channel.



The End

