# Halo and Sub-Halo Finding in Cosmological N-body Simulations

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## **Cosmological N-body Simulations**

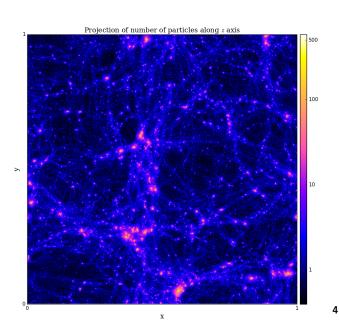
- ► N-body simulations are simulations of the motion of particles under the influence of physical forces.
- Focus on collisionless dark matter particles:
  - ► hypothetical type of matter
  - ► the only significant interaction between the particles is via gravity

## **Cosmological N-body Simulations**

- ► After some time, the particles will clump together. Such gravitationally bound objects are called *halos*.
- Halos themselves may contain self-bound objects, called subhalos.
- ► The identification of halos and subhalos is an important tool for problems concerning cosmic structure and its formation.
- ► Codes that perform this task are called *halo-finders*.

## **Cosmological N-body Simulations**

The results of a cosmological simulation of 128<sup>3</sup> dark matter particles at redshift z = 0with  $H_0 = 70.4$ and density parameters  $\Omega_m = 0.272$  and  $\Omega_{\Lambda} = 0.728$ . The box length corresponds to 88.8 Mpc.



## **Unbinding Particles**

- ▶ By convention, it is customary to treat all particles assigned to a halo as bound to it, even though from a strict energetic perspective they may not be.
- ► For subhalos, on the other hand, it is vital to identify and remove unbound particles:
  - Subhalos are located within a host halo and therefore expected to be contaminated by the host's particles
  - Usually subhalos contain far less particles than their hosts, so assigning particles to it without an unbinding procedure can influence its physical properties significantly.
- "Removing a particle" means here to assign it to the parent structure. This applies recursively to any level of substructure within substructure.

### Goals of this thesis

- RAMSES (Teyssier, R. 2002) is a N-body and hydrodynamical code that contains a clump finding algorithm, PHEW (Bleuler et al. 2015).
- ► Both PHEW and RAMSES are fully parallel and make use of the MPI library. PHEW works on-the-fly.
- ► The goal of this thesis is to implement a particle unbinding algorithm to work with PHEW that is also fully parallel and works on-the-fly.

### **PHEW**

- ► PHEW groups cells together by separating the mass density field along minima, thus dividing the density field into patches.
- ► The algorithm can be divided in four main steps:
  - ► segmentation
  - ► connectivity establishment
  - noise removal and
  - substructure merging

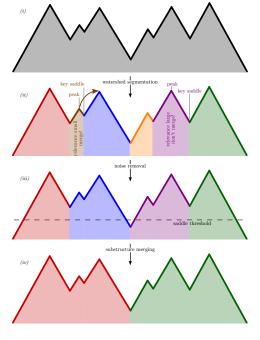
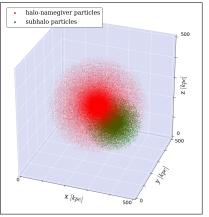


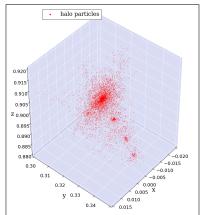
Image adapted from Bleuler et al. 2015.

### **Test Cases**

To demonstrate the effects of the particle unbinding, the following datasets will be used:

- dice-twobody-dataset: A highly idealistic structures where the effects can be seen and evaluated more easily, created using DICE (Perret 2016).
- cosmo-dataset: A halo from the previously shown cosmological simulation which is made up from 7030 particles.





The initial particle distribution of the dice-twobody dataset. A smaller halo (subhalo 1) made of 40'000 particles is nested within a bigger halo (halo-namegiver), which contains 200'000 particles.

cosmo-dataset: A halo as identified by PHEW of the previously shown cosmological simulation at redshift z=0.

## **Particle Unbinding**

In an isolated system in the centre of mass frame, each particle i can be assigned an energy  $E_i$ :

$$E_i = T_i + V_i = \frac{1}{2}m_i \cdot v_i^2 + m_i\phi(\vec{r}_i)$$

A particle is considered bound if:

$$E_i < 0 \quad \Leftrightarrow \quad v_i < \sqrt{-2 \cdot \phi(\vec{r}_i)}$$

## **Particle Unbinding**

The only considered potential  $\phi$  is the gravitational potential of the particles themselves. The potential is determined by the Poisson equation:

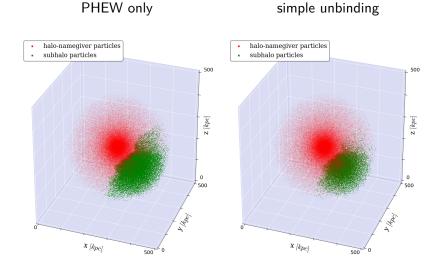
$$\Delta \phi = 4\pi G \rho$$

The spherically symmetric Poisson equation can be solved analytically for  $\phi$ :

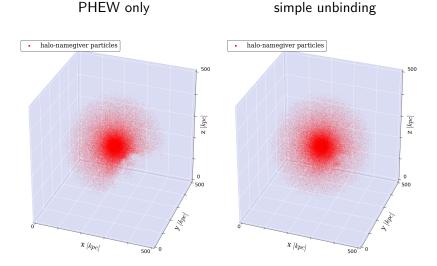
$$\phi(r_i) = -G \int_{r_i}^{r_{\text{max}}} \frac{M(<\tilde{r})}{\tilde{r}^2} d\tilde{r} - G \frac{M_{\text{tot}}}{r_{\text{max}}}$$

Where  $M(< r) \equiv \int\limits_0^r 4\pi \rho(\tilde{r}) \tilde{r}^2 \mathrm{d}\tilde{r}$  is the mass enclosed by a sphere of radius r such that the clump's total mass is enclosed by the radius  $r_{max}$ :  $M_{tot} = M(< r_{max})$  and G is the gravitational constant and  $\rho$  is the density.

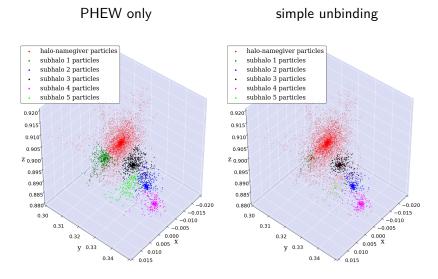
## Results: dice-twobody-dataset



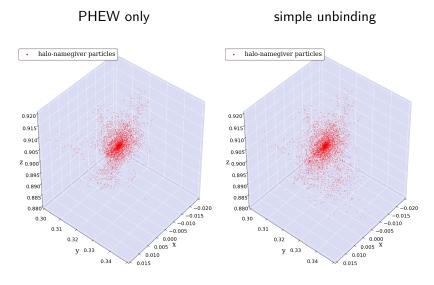
# Results: dice-twobody-dataset: halo-namegiver particles only



# Results: cosmo-dataset: halo-namegiver particles only



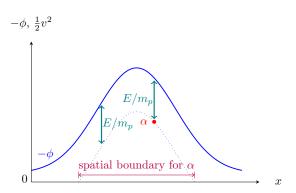
# Results: cosmo-dataset: halo-namegiver particles only



## **Accounting for Neighbouring Structures**

By construction, the identified subhalos are not isolated. This fact changes the situation significantly for the interpretation of what particles should be considered bound.

Consider first a particle  $\alpha$  in the potential of an isolated clump:

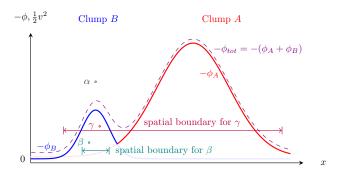


The spatial boundaries of its trajectory can be found by demanding energy conservation  $E/m_p=\frac{1}{2}v^2+\phi=const.$  by following the curve of constant total energy to the points where  $v^2=0$ .

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## **Accounting for Neighbouring Structures**

Now apply the same thoughts to an isolated halo that is made up from two clumps:



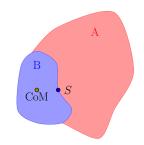
- $ightharpoonup \alpha$  is clearly not bound to the clump B.
- ightharpoonup eta will remain bound on an elliptic trajectory around the centre of mass.
- $ightharpoonup \gamma$  is energetically bound to the clump just like  $\beta$ , but because of clump A's neighbouring potential, the particle can leave the boundaries of clump B and wander off deep into clump A.

## **Accounting for Neighbouring Structures**

 $\Rightarrow$  Particles like  $\gamma$  shouldn't be considered bound.

The reason  $\gamma$  can wander off is because its boundary extends past the interface that connects the two clumps

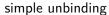
connects the two clumps ⇒ the condition for a particle to be *exclusively* bound must be that its trajectory must never reach that interface.

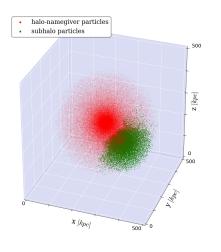


 $\Rightarrow$  Define S to be the point on the interface to the neighbouring structure(s) that is closest to B's centre of mass and  $\phi_S$  to be the potential of clump B at that point. Using the same argumentation as before, a particle can't reach S if

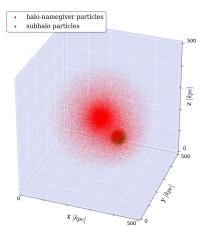
$$v<\sqrt{-2(\phi-\phi_S)}$$

## Results: dice-twobody-dataset

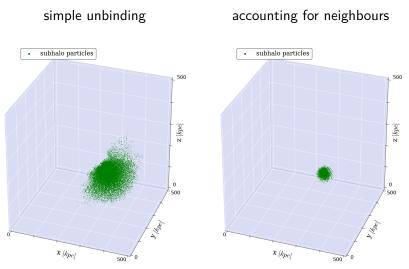




### accounting for neighbours



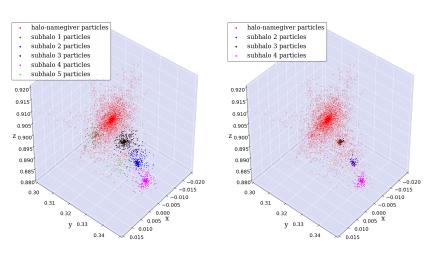
# Results: dice-twobody-dataset: subhalo particles only



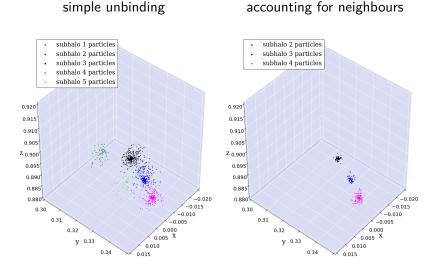
### Results: cosmo-dataset

### simple unbinding

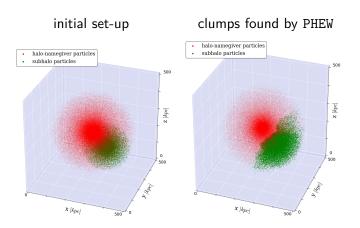
### accounting for neighbours



### Results: cosmo-dataset: subhalo particles only



## **Biased Clump Properties**



The identified clump properties will be biased:

- Missing particles: Subhalo is cut off
- ► Alien particles: Subhalo in contaminated by host's particles

## **Biased Clump Properties**

It seems likely that the clump properties after particle unbinding should be closer to the known ones, particularly so if only exclusively bound particles are considered.

⇒ recompute the clump properties after unbinding and use this updated information to go through the entire procedure again. Reiterate until the bulk velocity of each clump converges:

bulk velocity converged 
$$\Leftrightarrow \left| \frac{\textit{v}_{\textit{bulk},\textit{old}} - \textit{v}_{\textit{bulk},\textit{new}}}{\textit{v}_{\textit{bulk},\textit{old}}} \right| < \varepsilon$$

where  $\varepsilon$  is a user-defined convergence limit.

# Results: Converging of the bulk velocity for the dice-twobody-dataset

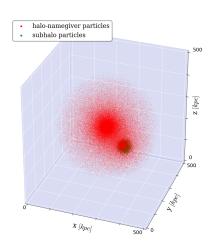
The deviation  $D_{orig} = \left| \frac{v_{bulk} - v_{orig}}{v_{orig}} \right|$  from the originally set bulk velocity to the computed bulk velocity for the subhalo of the dice-twobody-dataset in dependence of the convergence limit  $\varepsilon$ :

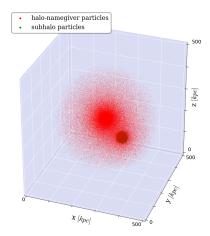
$\varepsilon$	niter	$D_{orig}$
0.5	2	0.2326
0.1	4	0.0419
0.01	7	0.0024
0.001	8	0.0014
0.0001	10	0.0009

The bulk velocity computation needed niter iterations to converge.

## Results: dice-twobody-dataset

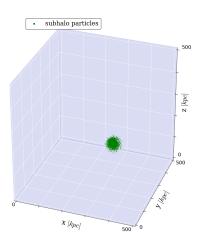
accounting for neighbours

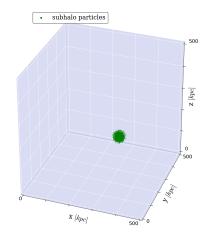




# Results: dice-twobody-dataset: halo-namegiver particles only

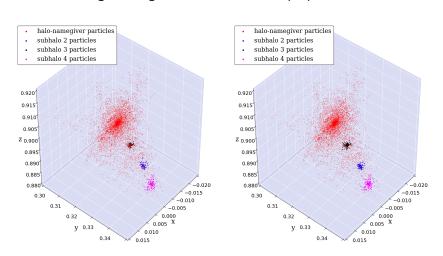
accounting for neighbours





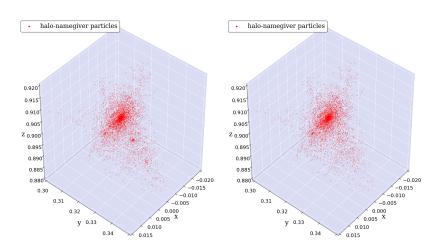
# Results: cosmo-dataset: halo-namegiver particles only

accounting for neighbours



# Results: cosmo-dataset: halo-namegiver particles only

accounting for neighbours



### References



Andreas Bleuler et al. "PHEW: a parallel segmentation algorithm for three-dimensional AMR datasets". In: *Computational Astrophysics and Cosmology* 2.1 (2015), pp. 1–16. ISSN: 2197-7909. DOI: 10.1186/s40668-015-0009-7. URL: http://dx.doi.org/10.1186/s40668-015-0009-7.



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